

## RESEARCH ARTICLE

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## Key Points:

- Gravity waves and downslope winds affect the transport and deposition of volcanic ash
- Meteorological-ash-dispersion modeling using WRF is able to resolve orographic process that leads to enhanced ash deposition
- The prescribed model plume height can be modified by orographic effects such as gravity waves near the volcano

## Supporting Information:

- Supporting Information S1
- Table S1

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## Orographic effects on the transport and deposition of volcanic ash: A case study of Mount Sakurajima, Japan

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**Abstract** Volcanic ash is a major atmospheric hazard that has a significant impact on local populations and international aviation. The topography surrounding a volcano affects the transport and deposition of volcanic ash, but these effects have not been studied in depth. Here we investigate orographic impacts on ash transport and deposition in the context of the Sakurajima volcano in Japan, using the chemistry-resolving version of the Weather Research and Forecasting model. Sakurajima is an ideal location for such a study because of the surrounding mountainous topography, frequent eruptions, and comprehensive observing network. At Sakurajima, numerical experiments reveal that across the 2–8 $\phi$  grain size range, the deposition of “medium-sized” ash (3–5 $\phi$ ) is most readily affected by orographic flows. The direct effects of resolving fine-scale orographic phenomena are counteracting: mountain-induced atmospheric gravity waves can keep ash afloat, while enhanced downslope winds in the lee of mountains (up to 50% stronger) can force the ash downward. Gravity waves and downslope winds were seen to have an effect along the dispersal path, in the vicinity of both the volcano and other mountains. Depending on the atmospheric conditions, resolving these orographic effects means that ash can be transported higher than the initial injection height (especially for ash finer than 2 $\phi$ ), shortly after the eruption (within 20 min) and close to the vent (within the first 10 km), effectively modifying the input plume height used in an ash dispersal model—an effect that should be taken into account when initializing simulations.

## 1. Introduction

Volcanic eruptions introduce large amounts of tephra (i.e., volcanic ash) into the atmosphere, creating a major environmental hazard. Even a small amount of volcanic ash (>2 mg m<sup>−3</sup>) [Langmann *et al.*, 2012] is a hazard for international aviation, capable of causing major disruptions [Bonadonna *et al.*, 2012]. After deposition, it can directly affect life, livelihoods (destroying crops and pastures), and infrastructure, for example, by causing roof and building collapses, damaging and disrupting electricity networks, clogging drainage systems, and contaminating water supplies [Wilson *et al.*, 2012; Hampton *et al.*, 2015]. In the long term, in cities near active volcanoes, ashfall can also have an indirect effect: exacerbating preexisting respiratory conditions [Baxter *et al.*, 1999; Horwell and Baxter, 2006; Hillman *et al.*, 2012], causing psychological stress, and imposing the economic burden of regular clean-up and maintenance of vulnerable infrastructure [Wilson *et al.*, 2015]. Accurate prediction of the transport and deposition of ash is therefore vitally important for hazard management and mitigation.

The transport and sedimentation of volcanic ash are complex processes, with the residence time and fall velocity of ash depending critically on its size [Carey and Sparks, 1986; Bonadonna *et al.*, 1998]. Coarse volcanic ash (>1 mm in diameter) is deposited quickly after an eruption and within a few tens of kilometers from the vent in a largely predictable manner [Carey and Sparks, 1986; Bonadonna *et al.*, 1998; Beckett *et al.*, 2015]. “Medium-sized” ash (diameters between 63  $\mu$ m and 1 mm) stays airborne for more time and is thus more influenced by the local and regional wind field. “Fine” ash (<63  $\mu$ m) stays airborne for longer periods of time: from days (as in the case of the 2010 Eyjafjallajökull, Iceland, eruption; [Webley *et al.*, 2012]) to years (e.g., the 1991 Mount Pinatubo, Philippines, eruption) [McCormick *et al.*, 1995]. The deposition of fine ash is still an open research topic, but deposition through aggregation (the joining of airborne ash particles) is widely accepted to exert first-order control [Bonadonna *et al.*, 2012; Brown *et al.*, 2012; Beckett *et al.*, 2015].

The atmospheric flow can control both the transport of ash particles via wind dispersal and the deposition. For example, cases of “forced deposition” of volcanic ash occur due to the formation of gravitational instabilities in the plume and/or orographically induced flows or waves, leading to strong downward advection. This can be achieved in at least two ways: via a localized increase in ash concentration, a process primarily controlled by plume dynamics [Manzella et al., 2015], and while the plume traverses complex topography and is impacted by orographic effects in the lee of mountains [Watt et al., 2015]. Precipitation, also a result of the particular atmospheric flow in a given day, has a direct effect on the deposition of volcanic ash and gases by hydrometeor scavenging, i.e., wet deposition [Kawaratani and Fujita, 1990; Witham et al., 2005].

An orographic flow is the result of atmospheric flow responding to isolated mountains or mountain ranges. Depending on a number of factors relating to the mountain topography, atmospheric stability, and wind speed, this interaction creates a variety of flow phenomena both on the windward side—such as flow stagnation and splitting—and in the lee—such as lee vortices, gravity (or *mountain*) waves, hydraulic jumps (abrupt upward displacement of air), and downslope winds [Smith, 1980; Smolarkiewicz and Rotunno, 1989; Durran, 1990; Elvidge et al., 2016]. The Froude number ( $Fr = U/(NH)$ , where  $U$  ( $\text{m s}^{-1}$ ) is the incoming flow speed,  $N$  ( $\text{s}^{-1}$ ) is the Brunt-Väisälä frequency, and  $H$  (m) is the height of an obstacle), is commonly used to characterize a flow based on expected orographic effects: broadly speaking a value less than 1 indicates a nonlinear “flow around” regime with flow splitting and lee vortices, while a value over 1 indicates a linear “flow over” regime with mountain wave activity [Smith, 1980]. Resolving orographic flows is essential for a large number of meteorological impacts, such as intense rainfall, floods, avalanches, and wildfire prediction [Meyers and Steenburgh, 2013]. The very few relevant studies on the influence of these processes on ash transport have drawn similar conclusions [Watt et al., 2015].

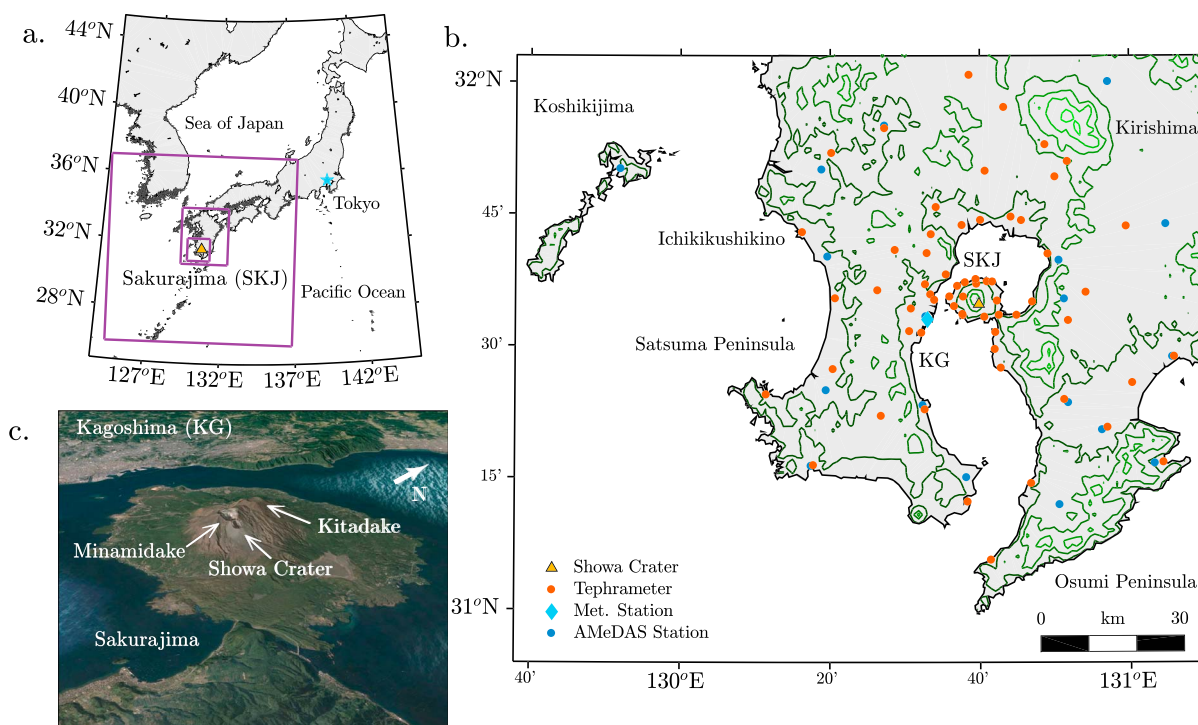
In this study we examine the impact of orographic effects on ash transport and deposition, in the context of the Sakurajima volcano in Japan using observational ashfall data and numerical modeling with a meteorological-ash-dispersal model. The localized nature of the dispersal around Sakurajima allows us to study the deposition using unprecedented horizontal and vertical resolution simulations. The paper is organized as follows. Initially, observational data showing the effect of orographic forcing on ash deposition will be examined (sections 2.1 and 2.2). The rest of the study focuses on the eruption of 18 August 2013 (introduced in section 2.3). The modeling setup (section 3) and a validation against observational data (section 4) are presented. Then the impacts of orographic effects are discussed in detail in section 5 and the conclusions in section 6.

## 2. Observed Features of Ashfall

### 2.1. The Sakurajima Volcano's Activity and Observations

Sakurajima (peak height 1117 m) is one of Japan's most active and closely monitored volcanoes, located on the island of Kyushu, southwestern Japan (Figure 1). The frequent activity of the volcano (currently ~80 events per month releasing approximately 500 kt of ash [Iguchi, 2016]) and a high population density have motivated several studies revealing the effect of the volcano on the well-being of the surrounding communities [e.g., Wakisaka et al., 1978; Yano et al., 1985, 1990; Uda et al., 1999; Samukawa et al., 2003; Horwell, 2007; Hillman et al., 2012]. Historically, the volcano had two main craters, Kitadake to the north (the earliest crater) and Minamidake to the south (Figure 1c). Activity in the volcano started again in 1955 mainly via the Minamidake crater. The latest eruptive phase, marked by vulcanian eruptions (i.e., moderate-sized, short-lived eruptive bursts) [Morrissey and Mastin, 2000] from the newly formed Showa crater to the south, started sporadically in 2008, with fairly constant activity since 2009 (see Figure S1 in the supporting information). Most eruptions are weak, but stronger eruptions with plume heights ( $H_p$ ) over 3 km (above ground level, agl) occur approximately every 5–6 weeks. In the latest eruptive phase, the tallest plume height reported has been 5 km agl in August 2013.

The Japan Meteorological Agency (JMA) provides forecasts of ashfall for eruptions with plume heights of over 2.5 km agl [Hasegawa et al., 2015], as well as a monthly report of activity [e.g., Japan Meteorological Agency (JMA), 2013]. These reports include a detailed analysis of major eruptions with simple maps of ashfall sightings communicated mainly via telephone to and by the JMA (for example, see the data included in Figure 5). Forecasts employ volcanic ash modeling via particle tracking using JMA regional (5 km horizontal grid resolution and 50 vertical levels) and local (2 km horizontal grid resolution and 60 vertical levels) model output [Shimbori et al., 2014]. The complex topography of the Kyushu island is known to have a strong effect on the atmospheric flow [Poulidis and Takemi, 2017]; however, these effects are not fully represented in the operational modeling currently carried out by the JMA.



**Figure 1.** (a) Location of the Sakurajima volcano (SKJ) in Japan. Purple squares indicate the three WRF model domains. (b) Map of the Kagoshima prefecture, with Kagoshima (KG), the Kirishima mountain, and other points of interest, as well as tephrameter (orange), and Automated Meteorological Data Acquisition System stations (AMeDAS; blue). Coastline data are shown with black contours, and topography contours are shown for every 300 m between 100 and 1600 m (dark to light green contours). Coastline is shown using the global, self-consistent, hierarchical, high-resolution shoreline data [Wessel and Smith, 1996], and topography contours are based on a digital elevation map (DEM) from the Advanced Spaceborne Thermal Emission and Reflection Radiometer (ASTER) mission [Yamaguchi et al., 1998]. (c) Three-dimensional rendering of Sakurajima and the surrounding topography from Google Earth.

Ashfall from the Sakurajima volcano has been monitored by the Kagoshima prefectural government since 1978 [Iguchi, 2016]. Tephrameters are used to measure ashfall at 62 locations (Figure 1b), daily for some stations within approximately 20 km from the vent and monthly otherwise. The JMA monitors eruptions by recording, for example, eruption time, plume height, and dispersal direction (for an example see Table S1). The eruption plume height is estimated using a combination of weather radar data from Fukuoka and Tanegashima (to the north and south, respectively) and visual data from a camera near the volcano [Shimbori et al., 2013a]. A meteorological station provides data on the vertical structure of the atmosphere (for example, temperature, wind, and stability) is in operation at Kagoshima [31.55°N/130.55°E, World Meteorological Organisation code: 47827] with rawinsondes launched twice daily [at 0900 and 2100 Japan standard time (JST = UTC + 9)].

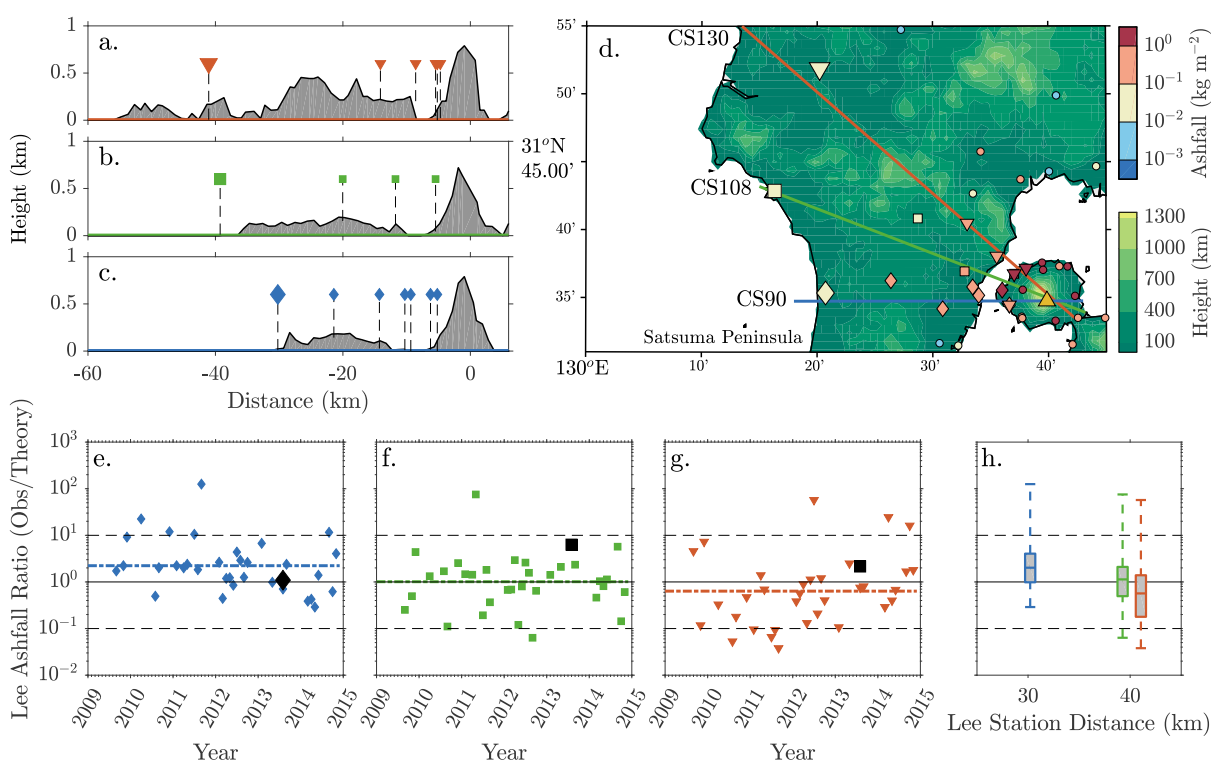
## 2.2. Observational Evidence of Orographic Forcing

The variation of ash thickness (and by extension ashfall amount) is a key diagnostic for estimating the erupted volume of ash [Daggitt et al., 2014]. In the case of the Sakurajima volcano, average monthly ashfall has been shown to decrease by an inverse power law with distance from the vent [Iguchi et al., 2013; Iguchi, 2016], following

$$w_a = w_M (D/D_0)^{-k}, \quad (1)$$

where  $w_a$  (in  $\text{kg m}^{-2}$ ) is the volcanic ash weight,  $D$  is the distance (in km),  $D_0$  is a reference distance used for nondimensionalization (1 km), and on a logarithmic plane,  $w_M$  (in  $\text{kg m}^{-2}$ ) is the intercept, and  $k$  is the slope [Iguchi, 2016]. The slope parameter (typically between 2 and 3 for Sakurajima) depends on the atmospheric conditions during the eruptions, such as the wind speed. The intercept coefficient,  $w_M$ , can be taken as an estimate for the ashfall at a distance of 1 km from the vent.

To investigate whether orographic effects have a significant impact, we analyzed monthly ash deposition observations. Wet deposition and orographically forced rainfall (i.e., rainfall triggered as the atmospheric flow



**Figure 2.** (a–c) Topography profiles at three cross sections west of the Sakurajima volcano. Markers indicate relevant tephrameter stations. Actual station height is shown by the dashed line. (d) August 2013 ashfall data over the Kagoshima prefecture for tephrameter groups along cross sections (CS) at different angles: CS90 (diamonds), CS108 (squares), CS130 (downward triangles), and others (circles). The Showa vent is marked with an upward triangle. (e–g) Ratio of observed to theoretical ashfall for the lee tephrameter stations along: CS90 (Figure 2e), CS108 (Figure 2f), and CS130 (Figure 2g). Black markers denote the August 2013 data points. Whisker diagrams of the lee ashfall ratio values versus lee station distance from the vent.

impinges on a mountain) [Houze, 2012] add a degree of complexity that makes a study of monthly data over-complicated. Thus, a subset of “dry” months are studied here. The criteria for selection are as follows: (i) months that featured at least one eruption with a plume height over 1 km dispersing toward the west or northwest and (ii) no rainfall observed west of the volcano for 6 h after the eruption (i.e., while the ash cloud is being advected over the land) in 90% of the eruptions of that month. For (i) JMA observations on the plume height and dispersal were used, while for (ii) hourly rainfall data from two rain gauge stations west of Sakurajima were used. This limited the study to 33 months (46% of the total period).

In the analysis we used data from three groups of stations: directly west (cross section 90, referred to here as CS90; diamond markers), west-northwest (CS108; square markers), and northwest (CS130; triangle markers) of the vent. The location of each station is shown in Figure 2d, while the topography profile along each cross section and the distance of relevant tephrameter stations is shown in Figures 2a–2c. The stations in each group are approximately aligned from the Showa crater so are expected to be affected by the same eruptions.

An increase in ashfall in the lee of a mountain is expected as a result of orographically forced deposition [Watt *et al.*, 2015]. For the case of Sakurajima we focused on the lee side of the Satsuma Peninsula, west of the volcano; specifically on the three most distal stations (shown with larger symbols in Figures 2a–2d). Orographic effects over the volcano not studied at this point, as ash from the large number of weak eruptions that occur every month, are expected to affect proximal ash deposits.

We compare the observed ashfall to that expected based on a power law (equation (1)). Using such a comparison highlights the departure from this simple fit that we associate with orographic effects (justified by our choice of dry months). For every month under study, power law fitting was applied on a number of tephrameter stations upwind of the lee station in order to derive a value for the slope coefficient ( $k$ ). This assumes that (i) the stations upwind of the lee stations adhere to this equation (which they do, the average coefficient of determination is over 0.9 for all groups) and (ii) the intercept coefficient over the stations are similar, which we can assume as they are on a roughly linear path from the vent. The ratios of the observed ashfall to

that expected with this derived power law slope ( $k$ ) value are shown in Figures 2e–2g. Values over 1 ( $10^0$  on the figure) indicate a departure from negative power law behavior, which is interpreted here as evidence of orographic processes modifying ash transport and deposition.

Overall orographic enhancement is most clearly noted for CS90; however, all groups show large scatter in the results. On average observed values are 23% larger (dash-dotted line) than the theoretically expected values. For most of the months under study ratio values are over 1, with the most extreme case showing a 100-fold increase. This effect is weaker for CS108 (overall average is close to 1), with almost half of the months suggesting increased deposition. Finally, the effect is further weakened for CS130 (average 66% of the theoretical value), suggesting no orographic enhancement on average. The lee ashfall ratio results are summarized with a box and whiskers plot in Figure 2h. In all cases the distribution shows orographic enhancement, i.e., a larger tail at higher ratios, with cases of sixtyfold enhancement even for groups CS108 and CS130.

Although cases of possible orographic enhancement exist for all groups, there is a large number of months when this process seems to be negligible. This can occur for a number of reasons. Probably, the most important one was mentioned previously: the effect of the topography would not appear explicitly in cases of weaker eruptions from which ash does not reach the lee stations. This would affect the ratio negatively and is expected to be exacerbated depending on the distance of the lee station, as seen here (Figure 2h). Atmospheric stability will also affect the results as atmospheric instability weakens gravity waves and thus affects deposition. Finally, in the case of CS130, the lee station is located in a valley between two mountains. Depending on the state of the atmosphere, this can also create interference patterns [e.g., Grubišić and Stiperski, 2009] that could weaken (or enhance) the effect. Further analysis is limited by the monthly nature of the data.

For the remainder of this study we focus on the month of August 2013 and specifically on the 18 August eruption as it has been the largest eruption to date of the current eruptive phase and the atmospheric structure on the day points toward the presence of orographic effects. Note that the August 2013 event is atypical meteorologically, leading to significant orographic enhancement along the CS108 line and no effect on the CS90 line (Figures 2e–2g). This is due to the meteorological conditions on the day.

### 2.3. The 18 August 2013 Eruption

On the 18 August 2013 Sakurajima erupted at 1631 JST (0731 UTC) with the plume height estimated between 5 and 7 km agl, based on the Kagoshima airport and JMA weather Doppler radars, respectively [JMA, 2013; Shimbori *et al.*, 2013b]. Ash was advected W-NW, with ashfall recorded as far as the Koshikijima islands, 90 km in the west (Figure 2c). The ash cloud can be seen moving toward a NW direction in the Multifunctional Transport Satellite (MTSAT-2) data (see Figure S2). The eruption was monitored by the Kagoshima airport Doppler radar in the north, a JMA weather Doppler radar on SKJ [Shimbori *et al.*, 2013b], and an X-band radar in the city of Tarumizu to the southeast of the volcano [Maki *et al.*, 2016].

This was the largest eruption since 2006, and even though it only lasted for a few minutes, it released an estimated  $334 \pm 33$  kt, equivalent to almost a month's worth of average ash emissions [see Iguchi, 2016]. We focus on this case for numerical simulations of the transport and deposition of the ash, in order to provide further insight on the orographic effects that appear to have affected the deposition.

## 3. Model Configuration and Experimental Design

Simulations of the 18 August 2013 eruption were carried out using the Weather and Research Forecasting (WRF) model [Skamarock *et al.*, 2008], version 3.6.1, coupled with “online” chemistry and aerosol calculations (WRF-chem) [Grell *et al.*, 2005]. WRF features a fully compressible, three-dimensional nonhydrostatic atmosphere, with the governing equations solved in flux form (Appendix A1).

Use of the WRF model has the benefit of online calculations: aerosol and meteorological calculations are carried out at the same time, eliminating the need for interpolation of the meteorological state between time steps and allowing for interactions between the ash cloud and the meteorological evolution. This, along with a high resolution sufficient to represent complex atmospheric flow phenomena that occur close to mountains and volcanoes [e.g., Minder *et al.*, 2013; Poulidis *et al.*, 2016], marks WRF as a suitable candidate for the study. Note that with the option used here, volcanic ash is treated as a passive scalar, i.e., complex chemical interactions are not included (see Appendix A2 for further details). To avoid confusion, we will be referring to the model as WRF to emphasize this fact—even though the chemistry module (WRF-chem) needs to be used to include volcanic ash.



### 3.1. Domain and Parametrization Options

For representing the complex topography of the Kyushu island nested domains were utilized (Figure 1a). The outermost domain (grid spacing of  $\Delta x = \Delta y = 12.5$  km,  $100 \times 100$  grid points) includes the islands of Kyushu and western Honshu. The intermediate domain ( $\Delta x = \Delta y = 2.5$  km,  $125 \times 150$  grid points) is centered over Kyushu. The innermost domain ( $\Delta x = \Delta y = 0.5$  km,  $300 \times 300$  grid points; Figure 1b) is centered over Sakurajima and includes the Kagoshima prefecture and the Koshikijima islands and is used for all figures shown. In the control simulations, referred to here as “normal” topography simulations, the Global 30 arc sec elevation (GTOPO30)—a 30″ ( $\approx 800$  m at this latitude) digital elevation map (DEM) data set—is used [Skamarock *et al.*, 2008], with a minimal amount of smoothing to ensure numerical stability of the model for these high-resolution experiments with very steep terrain [Zhong and Chow, 2013]. In order to study the impact of the topography and resolved orographic effects, additional “flat” topography simulations were carried out over a domain with almost zero topography (i.e., topographic height was divided by 1 km in all three domains, leading to a maximum topographic height of  $\sim 1$  m). There are 90 levels in the vertical with a vertical grid spacing of 50 m in the lowest 2 km. The model top is at approximately 20 km (50 hPa). The outermost domain boundaries are specified using the European Centre for Medium-Range Weather Forecasts Interim Re-Analysis data set [Dee *et al.*, 2011], which has been used in the past for similar ash dispersal studies [Macedonio *et al.*, 2016]. The boundaries for the inner domains use one-way nesting. The model is initialized at 0900 JST 18 August 2013 and run for 24 h. The first  $\sim 7$  h of the simulations are spent on model “spin-up” (approximately 3 h) and waiting for the eruption time (1631 JST). A suite of physical parametrizations (radiation, microphysics, surface layer, and fluxes) is used (see Table S3 for details). Sensitivity tests using different microphysics and planetary boundary layer schemes were carried out, but the results remained qualitatively unchanged.

Note that the 500 m grid size of the innermost domain cannot sufficiently resolve all turbulent motion [Bryan *et al.*, 2003; Takemi and Rotunno, 2003] and falls within the “turbulence gray zone” or “terra incognita”: turbulence is neither entirely parameterized nor resolved by a large eddy simulation scheme, which can lead to the neglect of potentially important terms in subfilter-scale turbulence models [Wyngaard, 2004]. Although this is not ideal, due to computational restraints, it is current practice in a number of mesoscale studies [e.g., Kirshbaum and Durran, 2004; Minder *et al.*, 2013; Cécé *et al.*, 2014; Kirshbaum and Fairman, 2014; Nugent *et al.*, 2014].

### 3.2. Source Parameter Specification

The ability of any Volcanic Ash Transport Model (VATM) to produce a realistic ash transport and deposition pattern hinges on the determination of appropriate source parameters for the volcano and the eruption [Mastin *et al.*, 2009; Bonadonna *et al.*, 2012]. While requirements vary with each model, WRF requires plume height (from the surface;  $H_p$ ), grain size distribution (GSD), eruptive mass rate ( $\dot{M}$ ; here  $1.11 \times 10^6$  kg  $s^{-1}$ ), eruption duration ( $d_E$ ; here 5 min), and the eruption time ( $t_E$ ; 1631 JST). Here most of these parameters are well constrained with data either provided from the JMA ( $H_p$ ,  $t_E$ ) or estimated by the Sakurajima Observatory of the Disaster Prevention Research Institute, Kyoto University [ $\dot{M}$ ,  $d_E$ ; Iguchi, 2016].

The plume height ( $H_p$ ) is known to be affected by atmospheric stability [Tupper *et al.*, 2009]. Here we use three different  $H_p$  values (3, 4, and 5 km agl) in order to test sensitivity of the vertical development of the plume height to the atmospheric structure as well as the overall impact in the ash deposition results.

As a default, WRF uses a simplified “umbrella” shape for the distribution of ash particles within the plume. This is appropriate for large (at least subplinian) eruptions (see Appendix A2 for details). In the case of a short vulcanian eruption, the plume is more akin to a thermal, rising in the atmosphere and disconnecting from the vent [Morrissey and Mastin, 2000]. The case studied here was in between these situations: not strong enough to be classified as a subplinian eruption (the total mass released is  $\approx 10^8$  kg, less than the typical threshold of  $10^{11}$  kg for subplinian eruptions) [Cioni *et al.*, 2000] but relatively long, lasting between 5 and 20 min and producing a sustained volcanic column. Note that it is difficult to calculate an exact value for the duration; here we have produced an estimate of 5 min based on seismic amplitude data on the 18 August [Iguchi, 2016]. Sensitivity tests carried out for a 20 min duration showed negligible differences in the results, and thus, the duration time of 5 min was chosen. The correct representation of the vertical distribution of ash is important for ash dispersal modeling and using a simplified profile is an approximation. However, recall that the main aim of this study is to examine orographic impacts on transport and deposition, rather than plume characteristics.

The GSD is the most difficult source parameter to constrain, as several samples at different distances from the vent are required to create a representative distribution [Bonadonna and Houghton, 2005]. In the case of Sakurajima this can be especially challenging due to the fact that explosive eruptions can occur along with passive degassing on a single day. Due to these difficulties, Mastin *et al.* [2009] assigned “default” source parameters for 1535 volcanoes (the majority of the Holocene volcanoes), including Sakurajima. Here we used the “Silicic, brief” profile in Mastin *et al.* [2009] which focuses on medium to light ash (between 2 and 7  $\phi$  the distribution is 0.09, 0.22, 0.23, 0.21, 0.18, and 0.07) and has a default duration of 6 min which is close to the eruption studied here. Note that  $\phi$  refers to the Krumbein scale:  $\phi = -\log_2(d/d_0^{-1})$ , where  $\phi$  (or “phi”) is the grain size,  $d$  is the particle size in millimeters, and  $d_0$  is a reference size (1 mm). Although the grain size is a dimensionless number,  $\phi$  is commonly used as a unit to avoid confusion. Finally, note that particle aggregation [Costa *et al.*, 2010] is not included in the simulations. At its current state WRF does not support any volcanic ash aggregation parameterizations; however, they are currently being implemented (M. Stuefer, personal communication, 2016). Note that the treatment of the volcanic ash, plume distribution details, and settling law used here is described in further detail in Appendix A.

## 4. Model Result Validation

### 4.1. Atmospheric Flow

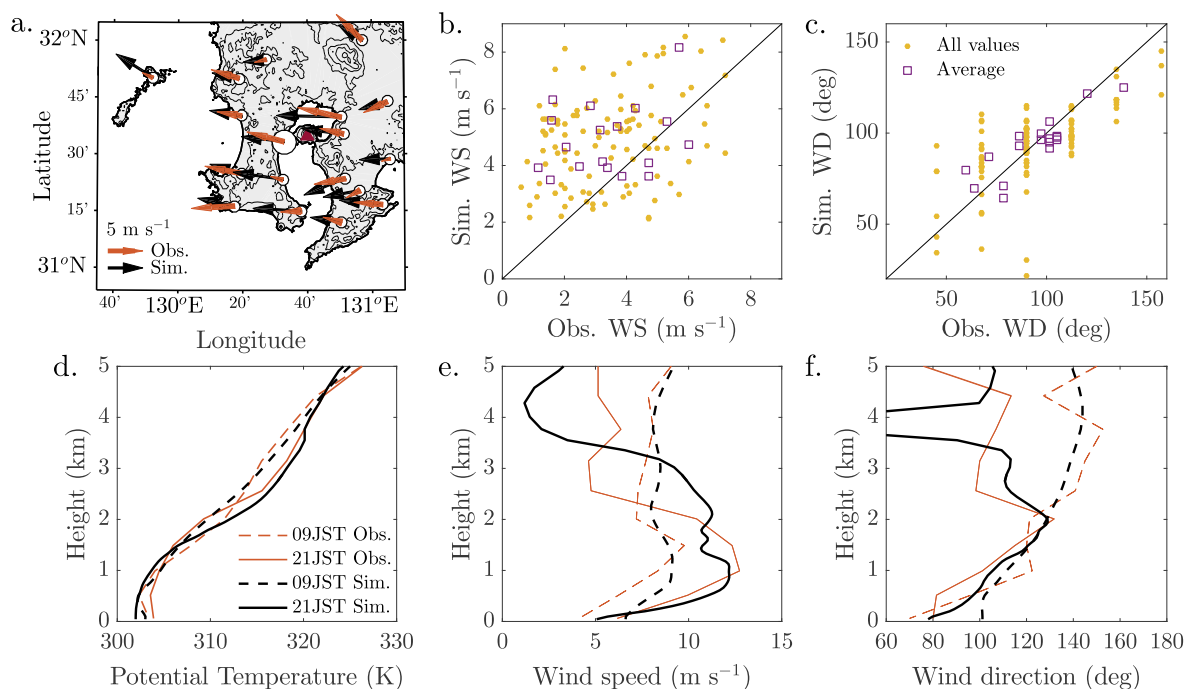
Surface observations (at 6–10 m) on the 18 August 2013 show atmospheric flow to the W-NW (Figure 3a). There are, however, several deviations from the average flow, indicating the influence of the topography; for example, the station to the north of the Kirishima mountain in the north of the domain (at approximately 32°N, 131°E), shows flow to the NW, suggesting flow splitting on the windward side of the mountain, and to a lesser degree at the Osumi Peninsula in the south (between 31°15′N–31°25′N and 130°40′–131°E), flow turns to the W-SW suggesting channeling.

In general, the simulated surface flow is in good agreement with Automated Meteorological Data Acquisition System (AMeDAS) observations, although for both instantaneous and 6 h average values, WRF tends to overestimate weak winds (Figure 3b): simulated surface wind is higher by a factor of 1.5–3 for  $U < 3 \text{ m s}^{-1}$  and corresponds better for stronger winds (Figure 3b). This is a known issue over complex topography and is thought to occur due to poorly resolved topographic features for atmospheric models, in general [Hanna and Yang, 2001], and WRF, in particular [Jiménez and Dudhia, 2012, 2013]. The root-mean-square error (RMSE) is  $2.44 \text{ m s}^{-1}$  for instantaneous and  $2.24 \text{ m s}^{-1}$  for hourly data. The wind direction correspondence is generally good: for instantaneous values the RMSE is  $18.95^\circ$ , while for 6 h averages this is further reduced to  $9.74^\circ$  (Figure 3c). In the case of both wind speed and wind direction RMSEs are within the expected range mentioned by Hanna and Yang [2001].

A comparison with rawinsonde observations over Kagoshima shows that the model is able to capture important characteristics of the vertical structure of the atmosphere (Figures 3d–3f). The potential temperature profiles match well, although the model produces a more stably stratified boundary layer (0–2 km; Figure 3d). The model captures the complex wind shear characteristics with some small deviations; specifically, the simulated decrease of wind speed occurs at a higher point (between 2 and 3 km instead of 2 km), and the model underestimates the minimum wind speed at a height of 4 km ( $1.2$  compared to  $5 \text{ m s}^{-1}$ ). The wind direction profiles (Figure 3f) are closely matched for the lowest 5 km with the exception of a relatively thin layer between 3.5 and 4.2 km, where there are very low wind speeds (another known WRF issue) [Jiménez and Dudhia, 2013].

The mountainous topography of the Kyushu island introduces significant orographic effects in the atmospheric flow as simulated in the WRF model (Figure 4). This is consistent with orographic flow theory as the Froude number for the case ranged between 1.1 and 1.35, indicating a flow over regime and the presence of gravity waves but with some flow splitting possible [Smith, 1980]. Note that Froude number was calculated upwind, for flow along the two main wind directions ( $108^\circ$  close to the surface and  $130^\circ$  at  $\approx 4 \text{ km}$ ), from the surface up to 2 km, and for an obstacle similar to Sakurajima (i.e., a height of  $\approx 1 \text{ km}$ ).

Close to the surface, strong E-SE winds introduce areas of ascent along the Osumi and Satsuma Peninsulas and on the windward sides of mountains, such as Sakurajima and Kirishima in the north (Figures 4a and 4c). Although the domain is mostly in a flow over regime, there are some cases of flow diversion, around the taller Kirishima mountain to the north (the peak height in the model being 1.5 km, leading to  $Fr = 0.79$ – $0.96$ ; Figure 4a) and to a lesser degree around Sakurajima (Figure 4c). On a vertical cross section, gravity waves can be clearly identified over and immediately downwind of Sakurajima (Figures 4e and 4f). These greatly alter



**Figure 3.** Comparison between simulated and observed meteorological data on 18 August 2013. (a) AMEDAS network (Obs.) and WRF (Sim.) averaged surface wind vectors. (b) Scatterplots of simulated wind speed against observational data using all data (yellow) and average values (purple). (c) As in Figure 3b but for wind direction. (d–f). Observed (orange) and simulated (black) vertical profiles of potential temperature (Figure 3d), wind speed (Figure 3e), and wind direction (Figure 3f). The vertical profiles shown are for the Kagoshima meteorological station (large marker size in Figure 3a).

the flow over the volcano introducing alternating areas of strong ascent (in red) and descent (in blue). Close to the surface (up to a height of 1 km), strong downslope winds develop (up to  $12.9 \text{ m s}^{-1}$ ), approximately 1.6 times larger than the 0–2 km layer average ( $7.8 \text{ m s}^{-1}$ ). The influence of the gravity waves is visible in the vertical structure of the wind fields up to 4–5 km in height to a distance of 40–50 km to the west at CS108 (Figure 4e) and continues beyond 60 km for CS130, enhanced by new activity over the Satsuma Peninsula mountains (between  $-20$  and  $-30$  km in Figure 4f). Note that the potential temperature lines (isentropes; brown lines Figures 4e and 4f) are parallel to the flow (as indicated by the wind vectors) as potential temperature is conserved for adiabatic motions. For this reason isentropes will subsequently be used to display flow lines.

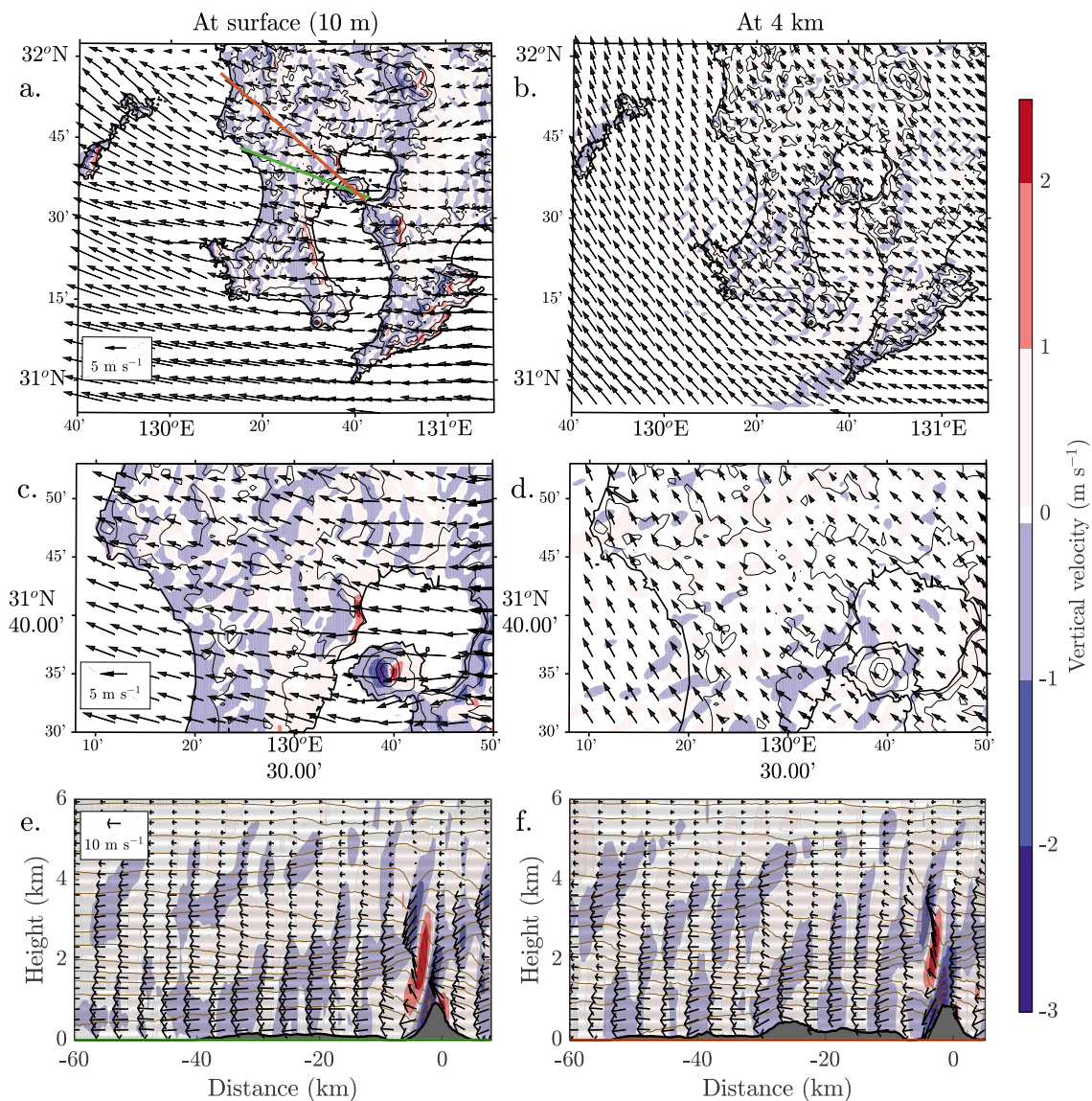
At a height of 4 km from the surface, the flow to the west of Sakurajima is southeasterly with considerably lower horizontal wind speeds to the NW of the volcano, especially over the land (Figure 4b). At this level, strong gravity wave activity can be identified in the southern part of the island by the alternating bands of blue and pink after the two ridge-like peninsulas and the Koshikijima islands.

Overall the WRF model is able to provide a reasonably realistic representation of the atmospheric flow on the day of the eruption. There is good agreement between simulated and observed wind direction. A number of orographic effects are resolved by the fine resolution of the simulations. Although the observational network around Sakurajima and the Kirishima mountain is not dense enough to capture detailed orographic flows in observations, the flow that was resolved by the model is as theoretically expected and there is good agreement between simulated and observed wind direction. On the other hand, comparison of simulated and observed wind speed shows larger disagreements; close to the surface low wind speeds are overestimated by the model, possibly due to subgrid topographic effects, a known issue for the model [Jiménez and Dudhia, 2012, 2013]. Finally, enhanced mountain wave activity due to the increased stability in the model will need to be kept in mind when discussing ashfall results.

#### 4.2. Ashfall

Ashfall patterns for the different plume heights ( $H_p$ ) and for both normal and flat simulations exhibit a number of similarities (Figure 5). All experiments show a NW distribution of ash, which is in agreement with JMA

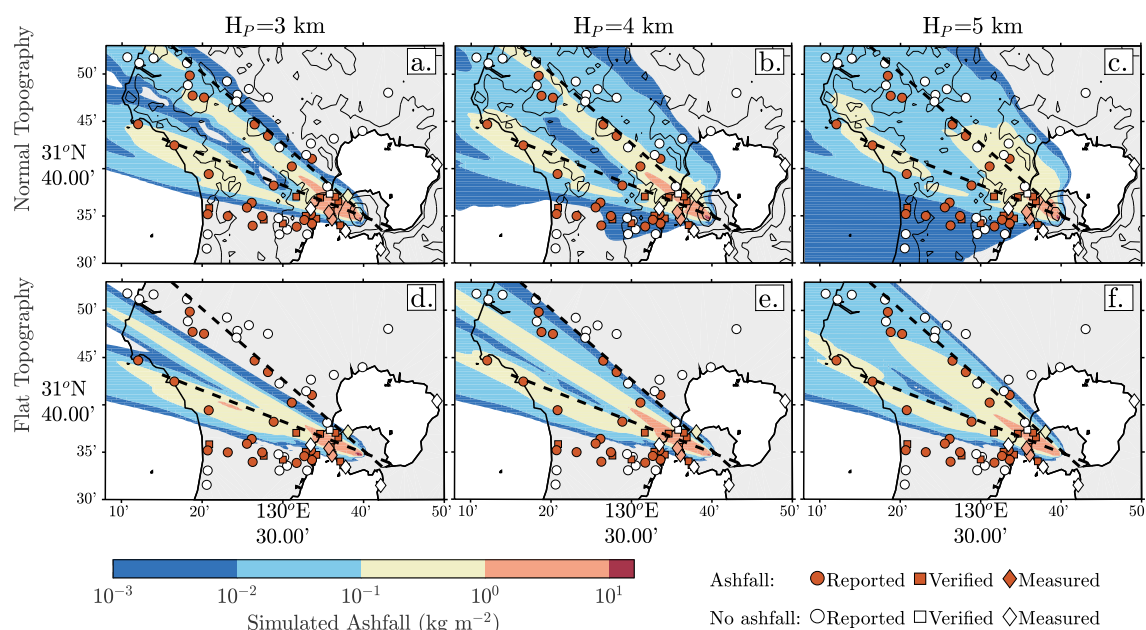




**Figure 4.** Simulated atmospheric flow for 18 August 2013. (a and b) Vertical velocity (shaded) and horizontal wind vectors at the surface (10 m) (Figure 4a), 4 km (Figure 4b) above the surface. (c and d) As in Figures 4a and 4b but focused over the west of Sakurajima. (e and f) Horizontal cross section of vertical velocity (shaded) with wind vectors and isentropes plotted every 2 K between 298 and 334 K. Cross-section lines are shown as lines in Figure 4a, green (CS108) for Figure 4e, and red (CS130) for Figure 4f. All results are averaged between 1600 and 2200 JST.

observations as well as radar and satellite data close to the volcano [Shimbori *et al.*, 2013b; Maki *et al.*, 2016]. Deposition occurs along two clearly defined branches in all cases. This is due to the shear in the vertical wind profile, with the two branches corresponding to the wind direction values close to the surface (lower branch; CS108) and aloft (upper branch; CS130). This is illustrated by the change in the patterns with plume height: lower plume heights lead to narrower deposition patterns (i.e., the two branches are closer together and the ashfall area is more limited) with a more pronounced lower branch and ash deposited more homogeneously along the path of the dispersal, while higher plume heights lead to wider deposition patterns with a gap over the Satsuma Peninsula mountains. Overall, results from the normal and flat topography experiments are similar. However, flattening the topography leads to narrower deposition patterns (note the deviation of the upper branch from the wind direction aloft) with more clearly defined edges.

In both types of simulations there are two maxima, one close to the volcano between Sakurajima and the northern part of Kagoshima and a secondary maximum on the lee side, over the Satsuma Peninsula (Figure 5).

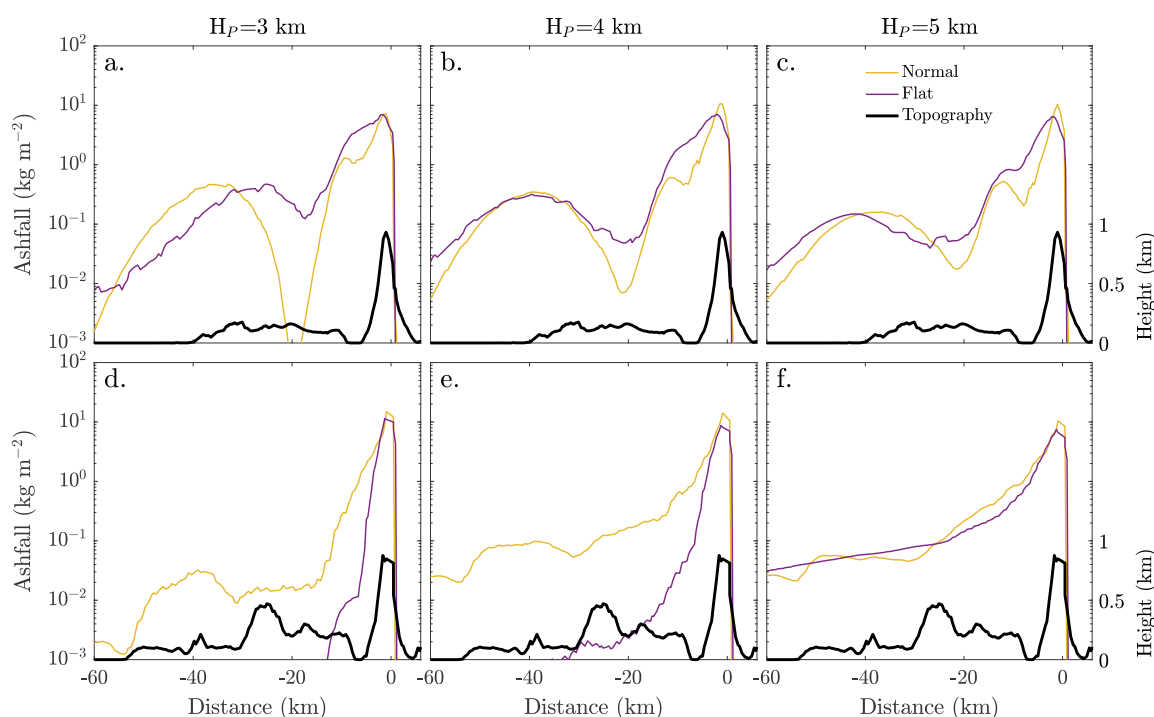


**Figure 5.** Simulated ashfall for (a–c) normal topography and (d–f) flat topography simulations. Results shown for different plume heights,  $H_p$ : 3 km (Figures 5a and 5d), 4 km (Figures 5b and 5e), and 5 km (Figures 5c and 5f). Ashfall observations are as reported from the JMA for the 18 August 2013. Simulated ashfall is shown at the end of the simulation (17 h after the eruption). Dashed lines mark cross sections CS108 and CS130.

Figure 6 shows that across the southern branch (Figures 6a–6c) the two maxima can be seen in both the normal and flat topography simulations. Field observations of downwind increases in ash deposition (termed secondary thickness maxima) are commonly inferred to arise from particle aggregation, enhancing the rate of ash settling [Carey and Sigurdsson, 1982; Brown et al., 2012; Watt et al., 2015]. Since no aggregation is included in our simulations, the model results suggest that departures from simple downwind patterns of decreasing deposition must be caused primarily by the topography, meteorology, and the depositional patterns of each variable ash bin. Although secondary maxima appear in both simulation types, the simulations that include topography produce a more pronounced secondary depositional maximum and shift this maximum more distally on the lee side of the Satsuma Peninsula.

Ashfall on the day was measured at a small number of locations near the volcano (indicated with diamond markers in Figure 5; also see Table S2). Although the number of the observation points is too small to conduct meaningful statistical analysis, they can be used for a rudimentary comparison with the model results. On Sakurajima, ashfall was measured at four points close to the coast: in the north ( $0.25 \text{ kg m}^{-2}$ ), west (approximately 1 and  $4 \text{ kg m}^{-2}$ ) and the south ( $0 \text{ kg m}^{-2}$ ). The observed value at the northern point is only captured well in the  $H_p = 4, 5 \text{ km}$  simulations with normal topography. In the  $H_p = 3 \text{ km}$  normal topography simulation, as well as in all the flat topography simulations, the depositional pattern is too narrow. The observed values at the points on the west of Sakurajima are captured better in the  $H_p = 3, 4 \text{ km}$  simulations, as the largest ashfall values shift to the north for the  $H_p = 5 \text{ km}$  simulations. At all of the points on Sakurajima ashfall is underestimated by the model in all simulations; this may be due to the lack of heavier ash that was not included in the GSD set in the model. Finally, over Kagoshima there is one point where recorded ashfall was explicitly  $0 \text{ kg m}^{-2}$ —in close proximity to locations where ashfall was observed. This (and other similar locations with mixed results in the JMA observational data) could be the result of microscale circulations within the city. Despite the high resolution, such effects are still subgrid scale and not resolved.

Along CS130, the normal topography simulations have a secondary depositional maximum in the lee of the Satsuma Peninsula mountain (at approximately  $-50$  to  $-25 \text{ km}$ ). This is clear in the  $H_p = 3 \text{ km}$  simulation (Figure 6d) but less clear with higher plume heights. The secondary depositional maximum in the simulation is associated with enhance downward deposition advection by downslope winds. In the  $H_p = 3 \text{ km}$  case the ash cloud is closer to the ground and thus more readily affected. This secondary maximum is not seen at all



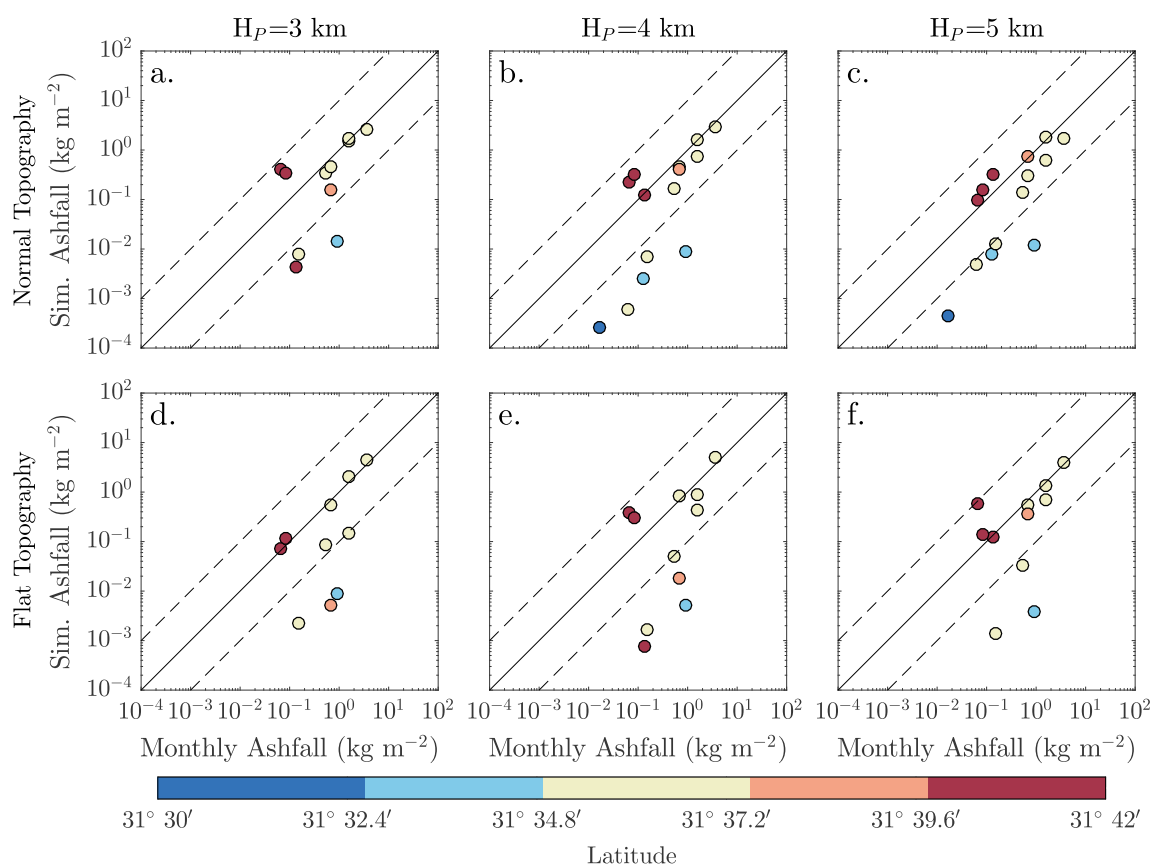
**Figure 6.** Ashfall distribution for the normal topography (yellow line) and flat topography (purple line) at cross sections shown in Figure 5. (a–c) CS108 and (d–f) CS130 for different plume heights. Topography (black line) shown in all panels to facilitate comparison. Note that the topography is not shown in a logarithmic scale.

in the flat topography simulations, even in cases where there is ashfall over the cross section (for example, Figures 6e and 6f).

As detailed depositional data are only available very close to the volcano, the simulated ashfall pattern away from the volcano can only be compared to the JMA report for the day [JMA, 2013]. As there is no specified lower limit for the JMA observations, here we use  $0.01 \text{ kg m}^{-2}$  for the calculations and took the position of each reported point as representative of the corresponding grid point in the model. This allowed us to estimate an “accuracy” value (i.e., the percentage of matching points) for each simulation. The overall accuracy of the simulations is between 59 and 62% for the normal topography simulations and 38–55% for the flat topography simulations. In the former, the best overall accuracy occurs for the  $H_p = 5$  km plume simulation which, although overestimating ashfall northwest of the volcano (i.e., north of the CS130 line), does include matches to the south and west. This is in agreement with the observations of column height. Similar conclusions can be drawn when comparing the simulated ash cloud to satellite data (see Figure S2).

A quantitative comparison of ashfall away from the volcano is also possible, against monthly observations (Figure 7). For August 18 2013 the event was decidedly the largest and most significant eruption (see Table S1), releasing half of the total ash released during the month, and so an order of magnitude comparison is justified. For most stations to the northwest of the volcano (latitudes over  $31^\circ 34.8'$ ) simulated ashfall is within an order of magnitude of that observed, i.e., for 50–56% for the normal topography simulations and 46–50% for the flat topography simulations. Ashfall west of the volcano is underestimated in the simulations.

Despite a relatively good qualitative agreement between the model results and the observations, there are a few common points of disagreement. In both the 4 and 5 km plume height cases, the model tends to overestimate ashfall northwest of the volcano. More importantly, a common point in all simulations is relatively low ashfall values west of volcano over an area with a number of JMA ashfall observations. As seen in the detailed meteorological analysis in section 4.1, both at the surface and at 4 km there is a considerable southerly influence in the wind field. This pushes the ashfall distribution toward the north and makes the simulated ash deposition to the west of the volcano unlikely. Passive emissions from the volcano can last for many hours



**Figure 7.** Comparison of simulated ashfall values against monthly tephrometer values for August 2013 for (a–c) normal topography and (d–f) flat topography simulations and different plume heights. Marker color indicates the latitude of the tephrometer station.

[Iguchi, 2016], and on the day of the eruption the volcano was continuously emitting ash for most of the day. It is possible that ash from these constant emissions and other small eruptions traveled west of the volcano. However, the possibility of model error cannot be ruled out, especially as misrepresentation of wind direction over complex topography is a known issue in the model [Jiménez and Dudhia, 2013].

Other possible sources of discrepancy in the results might stem from the current configuration of the WRF model with respect to volcanic ash. The GSD is known to be an important factor in the final ash deposition patterns [Bonadonna et al., 1998]. Here we assume a “medium” GSD which is an approximation to commonly known GSDs in eruptions of similar scale to the event under study here. Because of the uncertainties in GSD, such discrepancies in the results can be expected. As such results presented here should be viewed as a semiidealized case study rather than an attempt to provide an exact representation of the event.

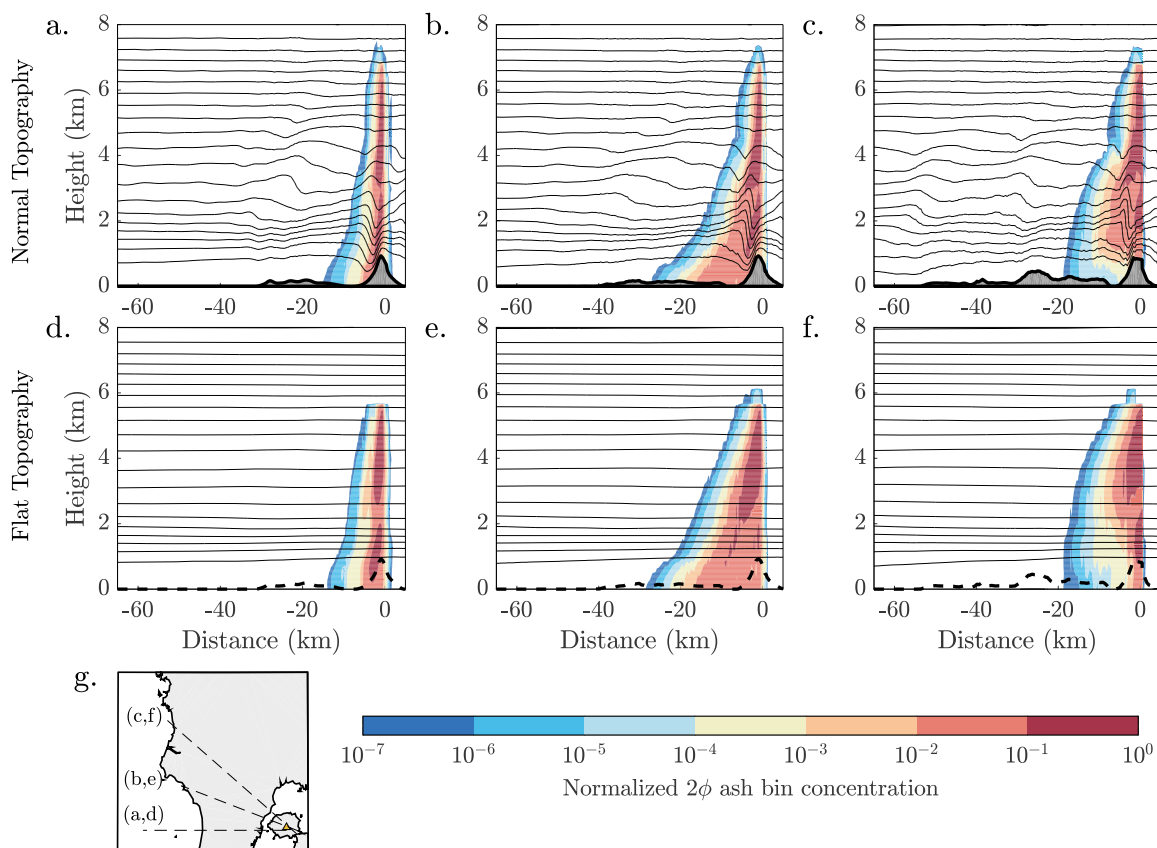
## 5. Orographic Processes Affecting Ash Dispersal and Deposition in Sakurajima

Comparison of the normal and flat topography simulations has shown that orography has a significant impact on ash deposition. The resolution of orographic effects has led to an overall increase in the fidelity of the simulation results. Here we will examine these effects in detail by looking at the dispersal of the ash plume. Note that all mentions of heights in this section are above sea level (asl) to avoid confusion when looking at the cross sections in Figures 8–10.

### 5.1. Orographic Effect Sensitivity by Grain Size

Figures 8–10 show vertical cross sections of the peak ash concentration (with respect to time) for the 2, 3, and 5 $\phi$  ash bins (examples of coarse, medium, and light ash, respectively) at three different angles passing through Sakurajima (see inset panels). Note that the concentration is normalized against the maximum





**Figure 8.** (g) Cross sections of maximum ash concentration (shaded) for the  $2\phi$  ash bin across the paths for (a–c) normal topography and (d–f) flat topography. Isentropes (6 h averaged) plotted as contours every 2 K between 298 and 334 K. For all simulations  $H_p = 5$  km.

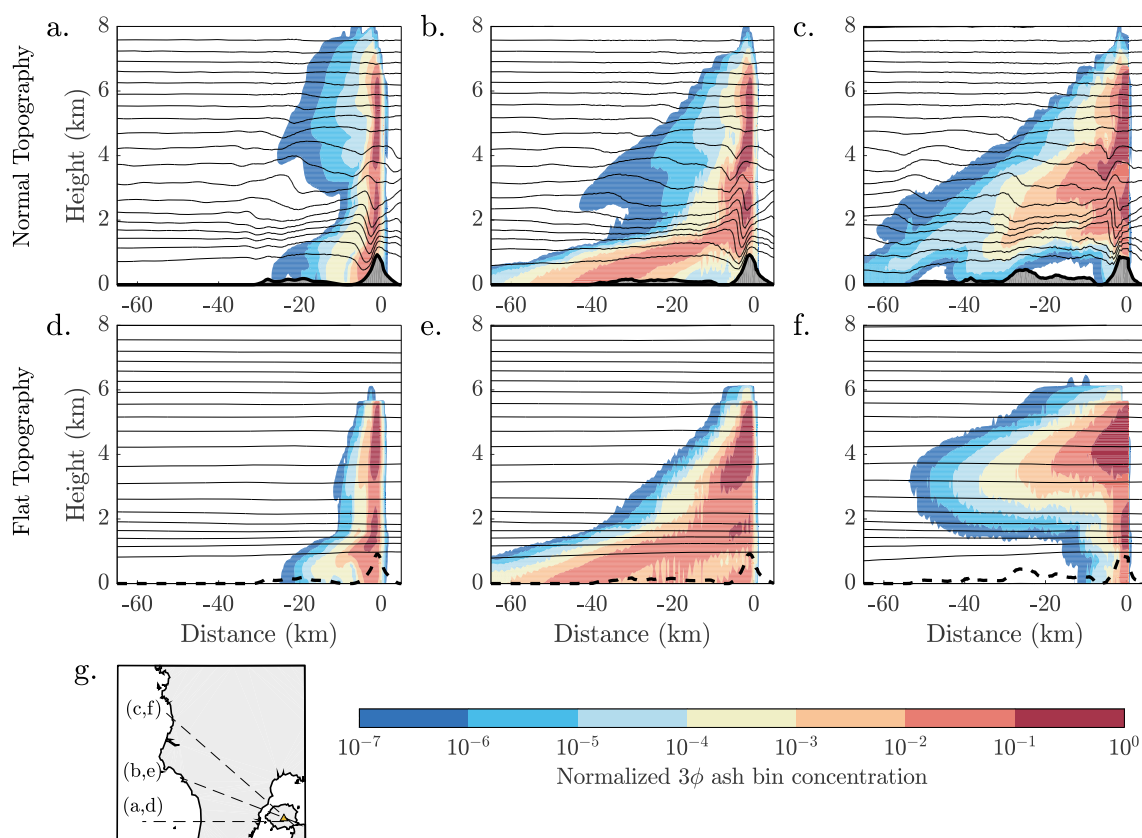
concentration value to allow an intercomparison. As the ash cloud moves from east to west without any points of flow reversal (right toward left in the figures), and the eruption duration was short with respect to changes in the wind field, the use of peak concentration provides a summary of the movement of the plume. The potential temperature lines (isentropes) are mean values for 6 h after the eruption and give an indication of the basic atmospheric flow during these 6 h. Due to the stability of the atmosphere, strong gravity wave activity is triggered in the lee of the volcano and in the lee of the Satsuma Peninsula (Figures 8a–8c).

The heaviest ash bin ( $2\phi$ ) is the least affected by orographic effects (Figure 8). The most noticeable difference can be seen for the main dispersal axes (Figures 8b and 8e, and 8c and 8f). Due to gravity wave activity over the volcano, ash is pushed farther up in the atmosphere, reaching over the originally prescribed volcanic plume height top of 6 km to 7.5 km asl. Close to the surface and directly west of the volcano, gravity waves keep ash afloat, reducing ash concentration and forcing ashfall further away.

The impacts of orographic effects are more clearly visible for the  $3\phi$  ash bin (Figure 9). It is clear that the ash distribution follows the gravity waves, being uplifted in the lee of the volcano and then subsiding farther downwind, so keeping ash afloat close to the volcano and focusing deposition over the lee of the Satsuma Peninsula acting along with the general forced descent due to the topography (compare Figures 9b and 9c to 9e and 9f). Close to the volcano, downslope winds have an opposing role to the gravity waves, as they push the lowest part of the plume downward, enhancing ashfall over the slopes (for example, Figure 9a). They can also have a more direct effect in the lee of other mountains as seen in Figure 9c where part of the ash cloud is pulled down toward the surface at  $-60$  km.

In the case of the  $5\phi$  ash bin (Figure 10), ash is largely deposited either near the volcano, mainly to the west, or outside the domain (Figures 10a and 10b). This is due to the combined effect of downslope winds over the volcano and increased mixing due to turbulent kinetic energy [Eckermann *et al.*, 2010]. Although not resolved in





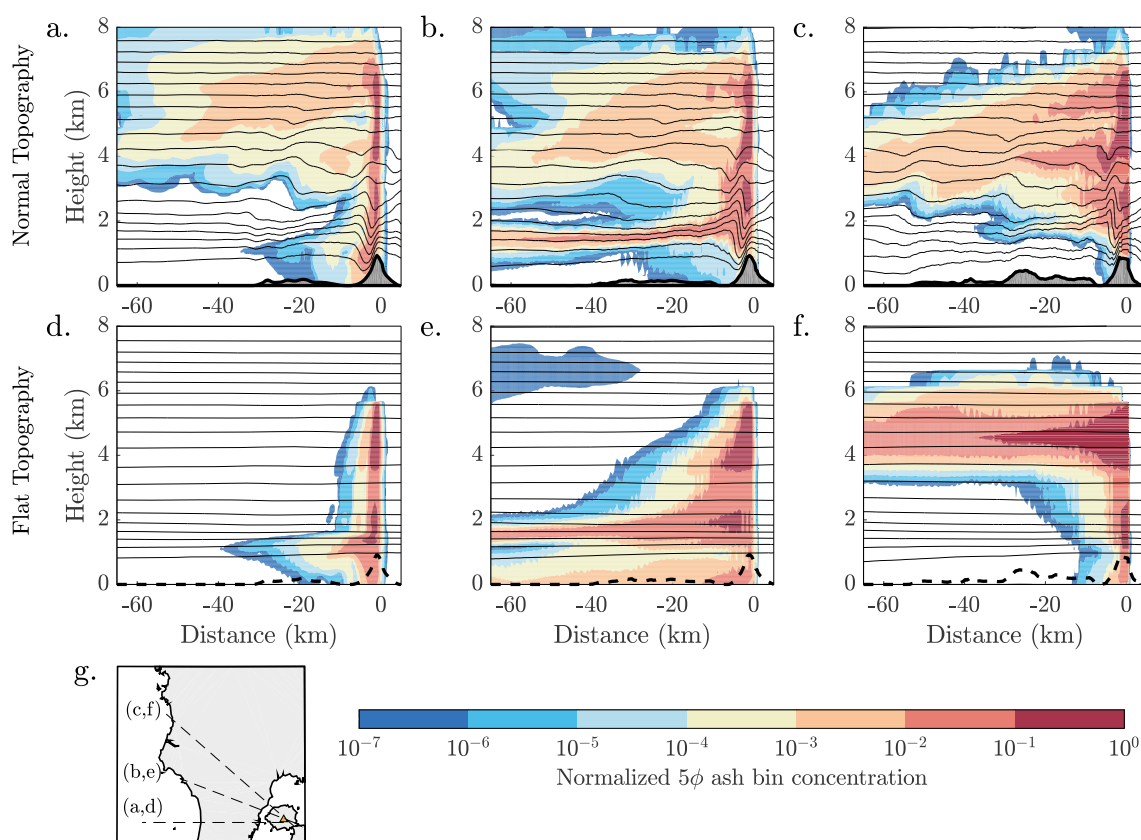
**Figure 9.** As in Figure 8 but for the  $3\phi$  ash bin.

the model, this orographically induced mixing close to the volcano could be enhancing gravitational instabilities that lead to proximal deposition of ash [e.g., Manzella et al., 2015]. Compared to the the flat experiments, there is less ashfall along the first NW direction (Figures 10b and 10e), as ash is trapped in a stratified layer between 1 and 2 km. Another direct effect of gravity wave activity is a large increase in ash concentration aloft (above 4 km asl) not present in the flat topography simulations. Note that lighter ash generally follows the same behavior as the  $5\phi$  ash bin.

## 5.2. Volcanic Plume Height

For the finer ash bins shown here (3 and  $5\phi$  ash bins), the ash is transported significantly higher than the initial plume injection height, due to upward advection by gravity waves. Within the first few minutes of the eruption in both the normal and flat simulations the top of the plume increases by  $\sim 1$  km due to mixing in the model. In the normal simulation this moves the top of the plume to approximately 7.5 km asl, which is close to the observed top [Shimbori et al., 2013b]. In addition, within the first 30–60 min and within the first 10–20 km from the vent the plume top rises by an additional 1–2 km up to 8.5 km asl. This is due to gravity wave activity occurring close to the volcano, pushing the air flow upward (for example, see between 5 and 8 km height, at  $-10$ – $0$  km in Figures 10b and 10c). This effect is less important for heavier particles that sediment quickly and are deposited within the first 30 min. Although not shown here, a similar increase is noted for the  $H_p = 3$  and 4 km simulations. This leads to a difference between the prescribed plume height ( $H_p$  in the present study) and the “effective” plume height (i.e., the maximum plume height simulated in the model over the volcano). An injection height of 6 km asl (i.e.,  $H_p = 5$  km) led to an increase of approximately 2.5 km, between 1 and 1.5 km over the highest radar-based estimates. An injection height of 5 km asl with a similar increase falls within the radar plume height estimate range, explaining the good performance of the lower plume heights.

In the eruption studied here, the increase in the plume height was exaggerated by the strong atmospheric stability. For similar atmospheric settings, the effect of gravity waves will be most acute in eruptions with plume heights less than 6–7 km. At the heights both gravity waves and downslope flows will appear significantly.



**Figure 10.** As in Figure 8 but for the  $5\phi$  ash bin.

## 6. Conclusions

Due to its frequent activity and comprehensive tephrameter network, the Sakurajima volcano provides an ideal natural laboratory for the study of orographic effects on ashfall. An analysis of tephrameter observations shows a clear enhancement of ash deduced to be orographically forced.

High-resolution meteorological-ash-dispersion modeling, capable of resolving complex orographic effects, has shown that the WRF model is able to produce realistic ashfall dispersal and deposition patterns. Medium-sized ash ( $3-5\phi$ ) was seen to be the most sensitive to orographic effects. There are counteracting processes with downslope winds and turbulent mixing enhancing ash deposition and gravity waves keeping ash afloat. Resolving orographic effects greatly enhances both horizontal and vertical diffusion of ash in the model. However, it leads to a potential ambiguity over the volcano: due to gravity wave activity over the volcano the prescribed “plume height” is elevated by up to approximately 2.5 km, an enhancement that could affect other low plume height eruptions.

Previous research has shown that eruptions and increased convection close to a volcano can affect gravity wave activity, while fine ash can be affected by atmospheric stability [Kanamori *et al.*, 1994; Tupper *et al.*, 2009; Poulidis *et al.*, 2016]. A similar effect was seen in the simulations here: resolving topographic effects can lead to increased ash concentration high in the atmosphere. Not resolving topographic effects could thus potentially lead to an underestimation of fine ash aloft with implications for aviation. The results here show high-resolution simulations are required to model and understand these complex processes. Note that the results from this study are based on one eruption. Further research for days with different synoptic conditions and in different topographic settings are needed to investigate the full extent of the impact of topographic effects on ash dispersal.

## Appendix A: WRF-chem as a Volcanic Ash Transport Model (VATM)

The WRF model is a versatile state-of-the-art numerical weather prediction model [Skamarock and Klemp, 2008] that has been used in a number of atmospheric studies, ranging from large eddy simulations of cloud dynamics [e.g., Kirshbaum and Durran, 2004] to regional climate change studies [e.g., Mearns et al., 2012]. The inclusion of the chemistry core [Grell et al., 2005] allows for the study of chemical pollutant and aerosol dispersal [e.g., Tie et al., 2007], as well as advanced aerosol-cloud interaction [Chapman et al., 2009]. Although use has been mainly constrained to light aerosols (particulate matter 10  $\mu\text{m}$  or less in diameter,  $\text{PM}_{10}$ , and  $\text{PM}_{2.5}$ ), it is possible to use built-in processes for the representation of volcanic ash [Stuefer et al., 2012]. Even though this has been used in a number of cases [e.g., Webley et al., 2012; Steensen et al., 2013], WRF-chem as a VATM is still at an early stage. For this reason some relevant characteristics of the model are included here for reference.

### A1. WRF-Chem Governing Equations

In WRF-chem the transport of all chemical species and aerosols is carried out online, at every time step, using the conservative (flux) form for the integration of prognostic equations [Grell et al., 2005]:

$$(\mu\phi)_t + \nabla(V\mu\phi) = 0. \quad (\text{A1})$$

In this equation  $\mu$  is the column mass of dry air,  $V$  is the velocity ( $u, v$ ), and  $\phi$  is a scalar mixing ratio. The model conserves mass and scalar mass exactly due to a finite volume formulation [Lin and Rood, 1996]. A Runge-Kutta third-order time scheme is coupled with a fifth-order evaluation of the horizontal flux divergence in the scalar conservation equation and a third-order evaluation of the vertical flux divergence [Wicker and Skamarock, 2002]. A positive-definite advection scheme is used for momentum, scalar, turbulent kinetic energy, and chemical species [Skamarock et al., 2008]. In the experiments here, a Rayleigh damping layer is used in the last 5 km before the model top [Klemp et al., 2008].

### A2. Volcanic Emissions and Plume

In WRF-chem, the eruption is not explicitly simulated; instead, a volcanic ash plume is inserted manually, in a predefined distribution, up to a specified height at a chosen time and sustained for the duration of the eruption. Poulidis et al. [2016] have shown that volcanic vent heating generates strong convective plumes which locally affect the atmospheric circulation. However, at present an eruption cannot be simulated in the WRF model. The default plume distribution used in WRF is the umbrella shape associated with large eruptions [Sparks et al., 1997], although this can be changed by the user. WRF uses a linear mass distribution beneath the umbrella—25% of the mass is contained in up to 73% of the eruptive column height, and a parabolic mass distribution with 75% of the entrained mass located on the top [Stuefer et al., 2012]:

$$M_a(h) = \begin{cases} N_L \frac{h}{U_b} & \text{if } h \leq U_b \\ N_U \left[ \frac{h-U_b}{H_p-U_b} - \left( \frac{h-U_b}{H_p-U_b} \right)^2 \right] & \text{if } U_b < h \leq H_p \\ 0 & \text{if } h > H_p \end{cases} \quad (\text{A2})$$

$$N_L = 0.25 \sum_{h=0}^{U_b} \left( \frac{h}{U_b} \right) M_T, \quad (\text{A3})$$

$$N_U = 0.75 \sum_{h=U_b}^{H_p} \left[ \frac{h-U_b}{H_p-U_b} - \left( \frac{h-U_b}{H_p-U_b} \right)^2 \right] M_T, \quad (\text{A4})$$

where  $M_a$  is the volcanic ash mass at height  $h$ ,  $N_L$ , and  $N_U$  are the mass beneath and included in the umbrella, respectively (also accounting for the normalization of the mass function),  $U_b$  is the umbrella bottom height,  $H_p$  is the plume height, and  $M_T$  is the total mass released.

There are 3 options for the representation of ash particles in the model, as 4 or 10 invariant tracer bins (grain size between 5–8 and –1–8, respectively), or 2 unspeciased aerosol bins ( $\text{PM}_{2.5}$  and  $\text{PM}_{10}$ ). In the experiments here we used the second option: Ash is inserted as an invariant tracer using 10 ash variable bins (grain sizes between –1 and 8  $\phi$  or diameters between 2000 and 4  $\mu\text{m}$ ) that are transported and settled. As no rainfall west of the volcano was observed on the day nor produced in the model, no wet deposition scheme was employed. The settling velocity uses an algorithm originally developed for the Goddard Chemistry Aerosol

Radiation and Transport model [Chin et al., 2002], based on the Stokes law, modified for volcanic ash radius and density. A Cunningham slip factor was also applied to account for noncontinuum effects for particles with diameters less than 15  $\mu\text{m}$  [Pruppacher and Klett, 1997].

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